

Solution to new sign problems with Hamiltonian Lattice Fermions

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- Can be written as

$$H = t \sum_{xy} c_x^\dagger M_{xy} c_y, \quad (2)$$

where

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- Particle-hole symmetry: $c_x \rightarrow \sigma_x c_x^\dagger$, $\sigma_x = (-1)^{x_1+x_2}$

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- The solution? Fermion bag approach.

The Naive Method

- We begin with writing $Z = \text{Tr} (e^{-\beta\epsilon})$ as

$$Z = \text{Tr} \left(e^{-\epsilon H} e^{-\epsilon H} e^{-\epsilon H} \dots e^{-\epsilon H} \right) \quad (4)$$

where there are N factors such that $N\epsilon = \beta$.

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$$= \int [d\phi] e^{-S[\phi]} \det M(\phi) \quad (7)$$

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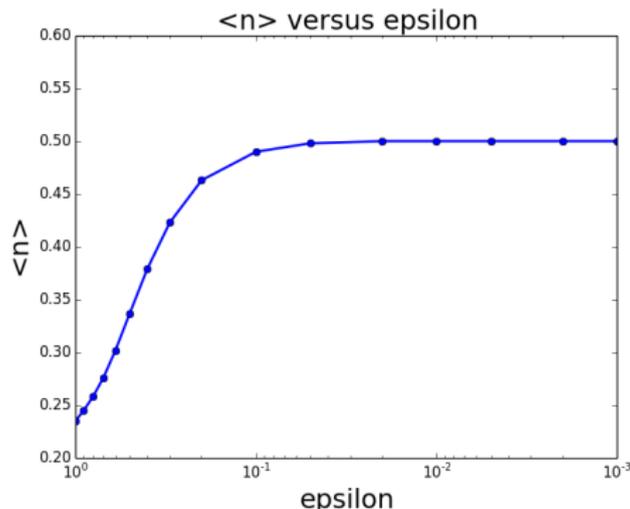
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$$\langle n_x \rangle = \frac{\int [d\bar{\psi} d\psi] e^{-S} \psi_x \bar{\psi}_x}{\int [d\bar{\psi} d\psi] e^{-S}}$$



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- We will see that, for a certain class of models, this expression may be written as determinants of matrices with some useful properties.

The Sign Problem in the Hamiltonian Approach

- Here we focus on a specific model involving staggered fermions:

$$H = t \sum_{x,y} c_x^\dagger M_{xy} c_y + \sum_{\langle x,y \rangle} \frac{V}{4} \left(n_x - \frac{1}{2} \right) \left(n_y - \frac{1}{2} \right) \quad (9)$$

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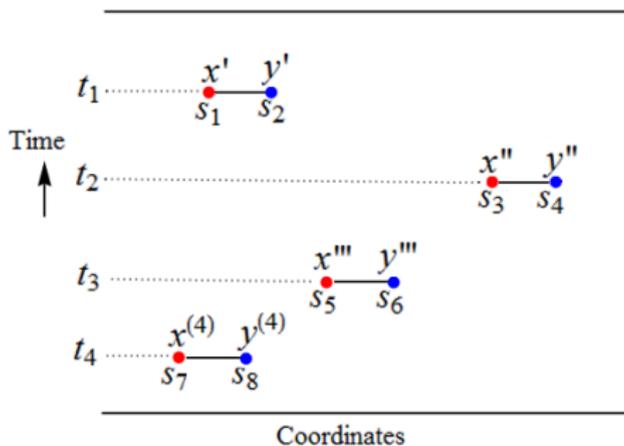
where $M'^T = -DM'D$, ($D_{xy} = \sigma_x \delta_{xy}$) 

The Partition Function

$$\begin{aligned}
 Z = Z_0 \sum_k \sum_{[b,s]} \int [dt] \left(-\frac{V}{4}\right)^k & \text{Tr} \left(e^{-(\beta-t_1)H_0} (s_{x'} n_{x'}^{s_{x'}}) (s_{y'} n_{y'}^{s_{y'}}) \right. \\
 e^{-(t_1-t_2)H_0} (s_{x''} n_{x''}^{s_{x''}}) (s_{y''} n_{y''}^{s_{y''}}) \dots & e^{-(t_{k-1}-t_k)H_0} (s_{x^{(k)}} n_{x^{(k)}}^{s_{x^{(k)}}}) (s_{y^{(k)}} n_{y^{(k)}}^{s_{y^{(k)}}}) e^{-t_k H_0} \left. \right)
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- The following identities hold: $a_{yx} = -\sigma_x a_{xy} \sigma_y$ and $d_{xx}[s] = -\frac{S_x}{2}$.

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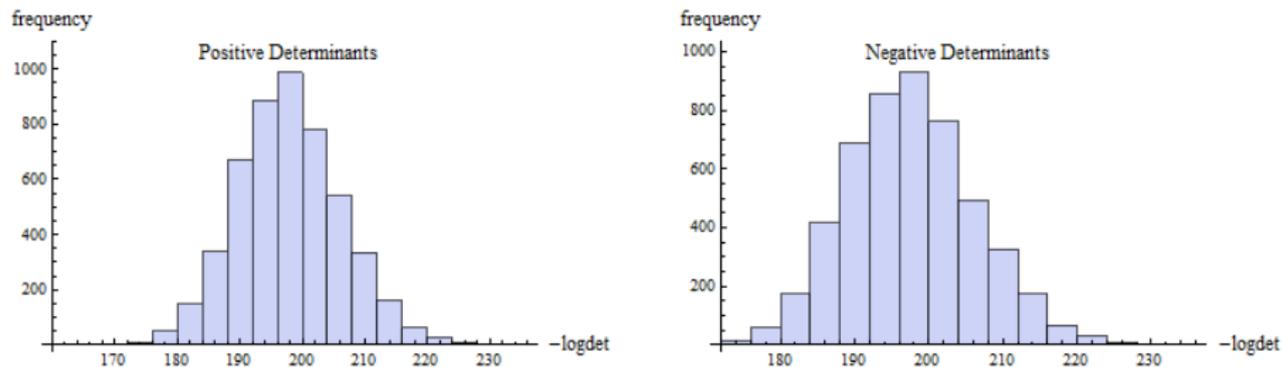


Figure: 10,000 determinants: 5004 were positive and 4996 were negative.

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- We first sum up the diagonal portion.

The Diagonal Sum

- We note that for the diagonal part:

$$\sum_{[s]} e^{-\bar{\psi} D_0([s]) \psi} = \prod_q \sum_{s_q=1,-1} \left(1 + \frac{s_q}{2} \bar{\psi}_q \psi_q \right) \quad (18)$$

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- Thus our partition function is now given by:

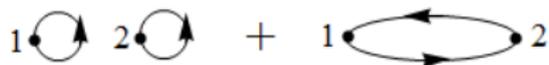
$$Z = \sum_{[b]} \int [dt] (-V)^k \text{Det}(A([b, t])) \quad (20)$$

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- In our sum of the $D_0 + A$ determinants, for every term of the form

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$$\begin{array}{c} \circlearrowright \\ \bullet \\ \circlearrowleft \end{array} \begin{array}{c} \circlearrowright \\ \bullet \\ \circlearrowleft \end{array} \dots \begin{array}{c} i \\ \circlearrowright \\ \bullet \\ \circlearrowleft \\ s_i = 1 \end{array} \dots \begin{array}{c} \circlearrowright \\ \bullet \\ \circlearrowleft \end{array}$$

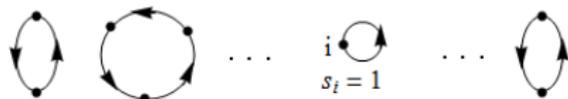
We have one with the form

Pictorial Proof

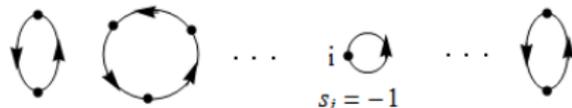
- Alternatively, we can see how this works using the pictorial representation of determinants. For example, a 2×2 determinant can be represented as:

$$1 \begin{array}{c} \circlearrowright \\ \circlearrowleft \end{array} 2 \begin{array}{c} \circlearrowright \\ \circlearrowleft \end{array} + 1 \begin{array}{c} \overbrace{\circlearrowright} \\ \underbrace{\circlearrowleft} \end{array} 2$$

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But are the determinants positive?

- $A([t])$ satisfies the relation $A^T = -\tilde{D}A\tilde{D}$, ($\tilde{D}_{xy} = \sigma_x \delta_{xy}$) so:

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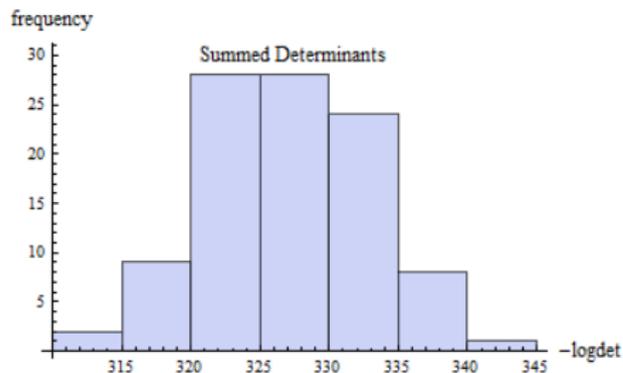
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- **We have solved the sign problem. (For repulsive model!)**

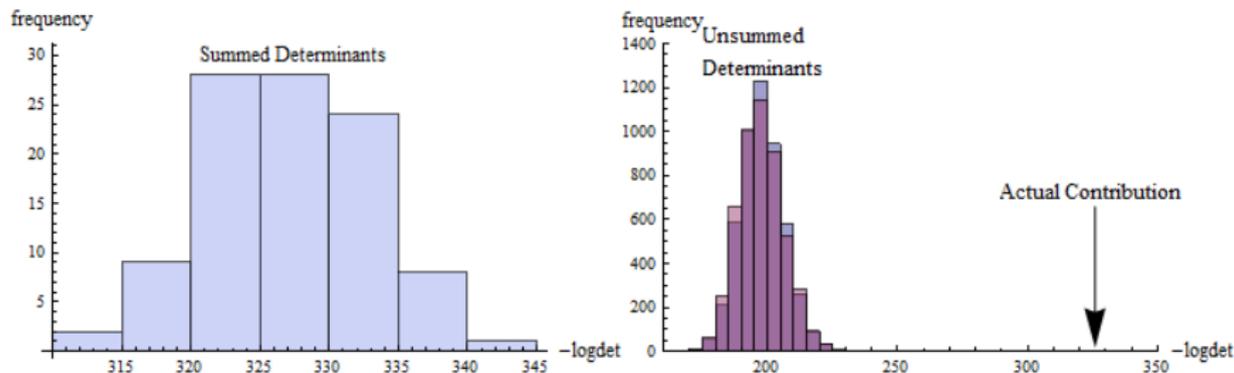
Some Example Determinants

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- Note that the probability of positive weight configurations is exponentially smaller, because the $-\log\det$ value is larger.



Conclusions and Future Work

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- Or we can add a staggered mass term that puts particles on the even sublattice and holes on the odd sublattice.
- Possible to study new quantum critical behavior.